

Cluster Compounds

Dimerization of Endohedral Fullerene in a Superatomic Crystal

Anastasia Voevodin,^[a] Laura Abella,^[b] Edison Castro,^[c] Daniel W. Paley,^[a, d] Luis M. Campos,^[a] Antonio Rodríguez-Forteza,^[b] Josep M. Poblet,^{*[b]} Luis Echegoyen,^{*[c]} and Xavier Roy^{*[a]}

Abstract: We describe a solid state material created from the reaction of $\text{Ni}_9\text{Te}_6(\text{PET}_3)_8$ and $\text{Lu}_3\text{N}@C_{80}$. The resulting superatomic crystal, $[\text{Ni}_{12}\text{Te}_{12}(\text{PET}_3)_8]_2[(\text{Lu}_3\text{N}@C_{80})_2]$, contains dimers of $\text{Lu}_3\text{N}@C_{80}$ that form upon reduction of the fullerene through a single C–C bond at the triple hexagon junctions. The encapsulated Lu_3N cluster displays an unprecedented orientation that is collinear and coplanar with the intercage carbon bond. Density functional theory calculations rationalize this unique bonding and relative orientation of the Lu_3N clusters. Our structural and theoretical results provide new insights into the effect that the M_3N cluster species has on the dimerization process of endohedral fullerenes.

The assembly of solid state materials from molecular building blocks offer significant benefits over traditional solid state reactions; the synthetic flexibility of the building blocks enables the development of functional materials with tunable properties. To this aim, fullerenes are attractive building blocks due to their exposed spherical π -surface capable of electronic coupling in all directions. Such electronic interactions in fullerene-based materials have enabled the emergence of remarkable collective properties such as ferromagnetism and superconductivity.^[1,2] Building on this foundation, our team has been developing a new class of solid state materials^[3] assembled from electronically and structurally complementary molecular clusters. These materials, which we term superatomic crystals,^[3–7] provide a bridge between traditional semiconductors,

molecular solids, and nanocrystal arrays by combining tunability and atomic precision. Fullerenes have been particularly useful to produce collective properties in superatomic crystals, including ferromagnetic ordering,^[8] coherent thermal transport^[9] and semiconducting behavior.^[3]

The synthetic flexibility of molecular clusters offers the possibility to create whole families of multifunctional materials by varying the constitution of the superatom building blocks. By contrast, the use of fullerenes has been, by and large, restricted to C_{60} and C_{70} . Endohedral fullerenes present the added benefit of varying the composition of the encapsulated guest while maintaining the advantageous properties of the carbon π -surface.^[10–12] These compounds have been explored as MRI contrast reagents, electron acceptors for photovoltaic cells, and single molecule magnets.^[13–15] Within the large family of endohedral fullerenes, metal nitride cluster fullerenes $\text{M}_3\text{N}@C_{80}$ stand out due to their compositional diversity and relatively high synthetic yields.^[15–17] Here we report on a new material, $[\text{Ni}_{12}\text{Te}_{12}(\text{PET}_3)_8]_2[(\text{Lu}_3\text{N}@C_{80})_2]$, which we discovered during our initial exploration of the reaction involving $\text{Ni}_9\text{Te}_6(\text{PET}_3)_8$ and $\text{Lu}_3\text{N}@C_{80}$. Using single crystal X-ray diffraction (SCXRD), we find that the anionic $\text{Lu}_3\text{N}@C_{80}$ fullerenes form dimers in this crystal. The triangular planar Lu_3N clusters inside the dimerized C_{80} cages are coplanar and collinear with the bridging C–C single bond, and point at each other. This observation contrasts with theoretical calculations^[18] and a recent experimental report of $[(\text{Sc}_3\text{N}@C_{80})_2]^{2-}$ dimers in which the clusters point away from each other.^[19] To understand this unusual orientation, we complement our experimental results with density functional theory (DFT) calculations. Our results chart a clear path to assembling novel superatomic crystals from endohedral fullerenes.

When compared to C_{60} or C_{70} , $\text{M}_3\text{N}@C_{80}$ have lower electron affinity.^[12,20,21] This presents an additional challenge for their assembly into superatomic crystals via charge transfer. To overcome this challenge, our initial plan was to react $\text{Lu}_3\text{N}@C_{80}$ with $\text{Ni}_9\text{Te}_6(\text{PET}_3)_8$, a building block with a high ionization energy.^[22]

Black crystals are obtained at the interface of two solutions containing $\text{Ni}_9\text{Te}_6(\text{PET}_3)_8$ and $\text{Lu}_3\text{N}@C_{80}$, dissolved in mixtures of quinoline and toluene, and kept at -35°C for seven days. SCXRD reveals that this solid is a 1:1 stoichiometric combination of a new cluster, $\text{Ni}_{12}\text{Te}_{12}(\text{PET}_3)_8$, and $\text{Lu}_3\text{N}@C_{80}$. Refinement of the crystallographic data indicates that $\text{Lu}_3\text{N}@C_{80}$ forms dimers and the stoichiometry of the compound is $[\text{Ni}_{12}\text{Te}_{12}(\text{PET}_3)_8]_2[(\text{Lu}_3\text{N}@C_{80})_2]$. Based on previous work on superatomic crystals,^[22] fullerene dimers^[17,23] and endohedral fullerene dimers,^[10,23] we can assign the following charges to the

[a] A. Voevodin, Dr. D. W. Paley, Prof. Dr. L. M. Campos, Prof. Dr. X. Roy
Department of Chemistry, Columbia University
New York, New York 10027 (USA)
E-mail: xr2111@columbia.edu

[b] L. Abella, Dr. A. Rodríguez-Forteza, Prof. Dr. J. M. Poblet
Departament de Química Física i Inorgànica
Universitat Rovira i Virgili, Marcel·lí Domingo 1
43007-Tarragona (Spain)
E-mail: josepmaria.poblet@urv.cat

[c] E. Castro, Prof. Dr. L. Echegoyen
Department of Chemistry, University of Texas at El Paso
El Paso, Texas 79968 (USA)
E-mail: echegoyen@utep.edu

[d] Dr. D. W. Paley
Columbia Nano Initiative, Columbia University
New York, New York 10027 (USA)

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building blocks: $[\text{Ni}_{12}\text{Te}_{12}(\text{PET}_3)_8]$ and the dimer $[(\text{Lu}_3\text{N}@C_{80})_2]$ bear 1+ and 2- charges, respectively. DFT calculations presented below agree with this assignment. Figure 1 presents the molecular structures of $\text{Ni}_{12}\text{Te}_{12}(\text{PET}_3)_8$ and the dianionic $[(\text{Lu}_3\text{N}@C_{80})_2]^{2-}$ dimer.

The cluster $\text{Ni}_9\text{Te}_6(\text{PET}_3)_8$ has been shown to reorganize in solution,^[24] and under our reaction conditions, it produces $\text{Ni}_{12}\text{Te}_{12}(\text{PET}_3)_8$ (Figure 1 a). The structure of $\text{Ni}_{12}\text{Te}_{12}(\text{PET}_3)_8$ can be

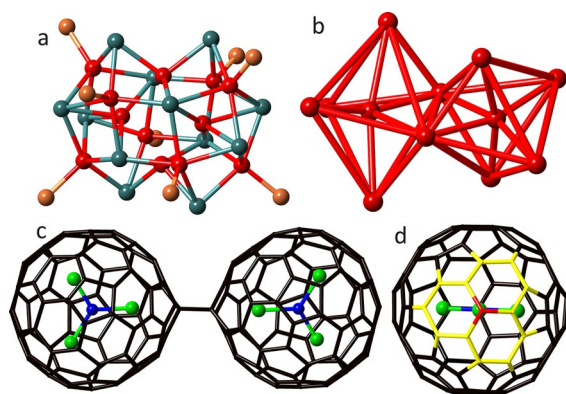


Figure 1. SCXRD molecular structure of (a) $\text{Ni}_{12}\text{Te}_{12}(\text{PET}_3)_8$. (b) Structure of the Ni framework with distorted octahedra sharing a vertex. The Ni–Ni bonds are added to highlight the shape of the octahedra. (c) $[(\text{Lu}_3\text{N}@C_{80})_2]$ dimer, and (d) $\text{Lu}_3\text{N}@C_{80}$ highlighting the THJ (yellow) of the inter-cage bonded carbon (red). Color code: C, black; N, blue; Lu, green; Ni, red; Te, teal and P, orange. Ethyl groups on the phosphines are removed from (a) to clarify the view.

viewed as two distorted Ni_6 octahedra, fused through a shared vertex. Each Ni_6 octahedron contains one interstitial Ni atom. The distances separating these interstitial Ni atoms and the Ni atoms at the corner of the octahedra range from 2.40 to 2.60 Å. For comparison, in metallic nickel the $\text{Ni}\cdots\text{Ni}$ distance is 2.49 Å. The other $\text{Ni}\cdots\text{Ni}$ distances are longer than 2.71 Å. The Te atoms adopt two binding modes: six Te bridge three Ni atoms and the other six Te bridge four Ni atoms. Phosphine ligands coordinate eight of the ten surface Ni atoms. The other two surface Ni atoms are bonded to four Te in a distorted tetrahedral geometry. A nickel selenide cluster with an analogous core composition, $\text{Ni}_{12}\text{Se}_{12}(\text{PET}_3)_6$, was previously reported by Fenske and Ohmer^[25] but the trioctahedral core of this compound is entirely different from the distorted structure of $\text{Ni}_{12}\text{Te}_{12}(\text{PET}_3)_8$ reported here.

The interesting result is that $\text{Lu}_3\text{N}@C_{80}$ is dimerized in the solid state, presumably as a consequence of the electron transfer from the electron-rich superatom (Figure 1c). Single-bonded dimers have been observed for reduced C_{60} and C_{70} , and suggested both theoretically^[18] and spectroscopically^[26] for endohedral fullerenes. Konarev and co-workers recently reported the first crystallographic evidence of dimerization of $[\text{Sc}_3\text{N}@C_{80}]^-$ upon reduction with sodium fluorenone ketyl.^[19] The $\text{Sc}_3\text{N}@C_{80}$ dimer structure features significant disorder of the cages, the inter-cage C–C bond and the cluster but the orientation of the Sc_3N cluster with respect to the inter-cage C–C bond is clear and in agreement with theory: the triangular

planar Sc_3N clusters are close to collinear with the bridging C–C single bond, and point away from each other.

The fully ordered structure for the $[(\text{Lu}_3\text{N}@C_{80})_2]^{2-}$ dimer differs significantly from that of $[(\text{Sc}_3\text{N}@C_{80})_2]^{2-}$.^[19] While the inter-cage C–C bond length (1.66(6) Å) for $[(\text{Lu}_3\text{N}@C_{80})_2]^{2-}$ is comparable to that for $[(\text{Sc}_3\text{N}@C_{80})_2]^{2-}$, it selectively links the hexagon-hexagon-hexagon junctions (THJ) of neighboring C_{80} cages (Figure 1d). This contrasts with the $[(\text{Sc}_3\text{N}@C_{80})_2]^{2-}$ structure, in which the fullerene dimer is disordered over three positions including a mixture of THJ and PHHJ dimers. More remarkably, one N–Lu bond for each Lu_3N cluster is perfectly collinear with the inter-cage C–C bond and points directly at the other Lu_3N cluster. While the C_{80} cages are fully ordered at 100 K, the Lu_3N clusters are disordered over three rotational orientations around the axis passing through the inter-cage C–C bond. The distance between the central N atom and the Lu atom closest to the inter-cage C–C bond (2.06(2) Å) is slightly elongated, when compared to the other Lu–N bonds (2.009(19)–2.03(3) Å). The Lu–N–Lu angles are close to the ideal 120° , ranging from $117.4(9)^\circ$ to $122.4(9)^\circ$ and the Lu_3N cluster is almost completely flat, with the central N atom protruding from the trimetallic plane by at most 0.039(14) Å.

Figure 2a shows the extended packing of $[\text{Ni}_{12}\text{Te}_{12}(\text{PET}_3)_8]_2[(\text{Lu}_3\text{N}@C_{80})_2]$, which can be visualized as the superatomic structural analogue of the binary atomic compound rubidium peroxide, Rb_2O_2 . Figure 2b compares both structures. We present schematic views of the superstructure in which a dummy atom is positioned at the center of each building block (blue represents $\text{Lu}_3\text{N}@C_{80}$ and red represents $\text{Ni}_{12}\text{Te}_{12}(\text{PET}_3)_8$). As with the peroxide dianion $[\text{O}_2]^{2-}$, pairs of blue atoms are linked together to represent the $[(\text{Lu}_3\text{N}@C_{80})_2]^{2-}$ dimers. The packing structures of Rb_2O_2 is presented looking down all three crystallographic axes, along with views showing the same orientations for the superatomic crystal. The superstructure of $[\text{Ni}_{12}\text{Te}_{12}(\text{PET}_3)_8]_2[(\text{Lu}_3\text{N}@C_{80})_2]$ presents a small distortion of the idealized Rb_2O_2 packing resulting from a tilt of the dimer with respect to the b-axis.

The relative orientation of the Lu_3N cluster within the dimer is unexpected. In fact, the question of the relative orientation of M_3N cluster upon exohedral functionalization has received little attention in the literature. To study this question, we performed DFT calculations for a series of $\text{M}_3\text{N}@C_{80}$ (with $\text{M} = \text{Lu}, \text{Sc},$ and Y). Table S3 contains the computed energy of each inter-cluster orientation upon exohedral dimerization through THJ and PHHJ junctions. Figures 3a and 3b illustrate the two possible orientations for the M_3N cluster. At the PBE/TZ2P level, the orientation in which the Sc_3N clusters are collinear, coplanar, and pointing at each other in the THJ dimer (orientation 1, Figure 3a) is energetically disfavored by more than 10 kcal mol^{-1} with respect to the opposite orientation in which the clusters point away from one another (orientation 2, Figure 3b). By contrast, orientation 1 is strongly favored by $4.0 \text{ kcal mol}^{-1}$ in the case of the Y_3N cluster. Lu_3N is the intermediate case as Lu sits between Sc and Y in terms of size and electronegativity. DFT calculations indicate that orientation 1, which is observed experimentally, is only slightly favored by $1.3 \text{ kcal mol}^{-1}$ over orientation 2.

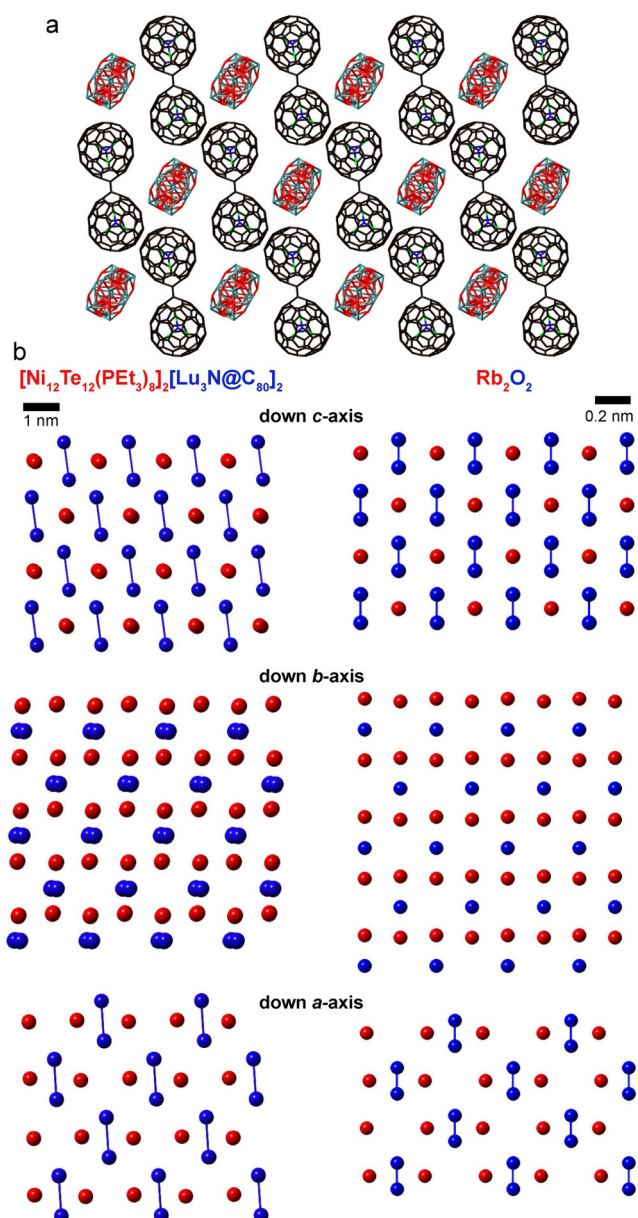


Figure 2. (a) Crystal packing of $[\text{Ni}_{12}\text{Te}_{12}(\text{PEt}_3)_8]_2[\text{Lu}_3\text{N}@C_{80}]_2$. Color code: C, black; N, blue; Lu, green; Ni, red; and Te, teal. The phosphines on the cluster and three quinoline molecules per formula unit are removed to clarify the view. (b) Schematic views comparing the packing of $[\text{Ni}_{12}\text{Te}_{12}(\text{PEt}_3)_8]_2[\text{Lu}_3\text{N}@C_{80}]_2$ and Rb_2O_2 . $\text{Ni}_{12}\text{Te}_{12}(\text{PEt}_3)_8$ and $\text{Lu}_3\text{N}@C_{80}$ are represented by blue and red dummy atoms positioned at the center of each building block, respectively. Pairs of blue atoms are linked together to denote the $[(\text{Lu}_3\text{N}@C_{80})_2]^{2-}$ dimer and the peroxide dianion $[\text{O}_2]^{2-}$.

We compare the energy differences for the various types of intercage bonding (i.e. THJ–THJ, PHHJ–PHHJ and THJ–PHHJ dimers shown in Figure 3) to understand why, unlike what has been reported for $[(\text{Sc}_3\text{N}@C_{80})_2]^{2-}$, the $[(\text{Lu}_3\text{N}@C_{80})_2]^{2-}$ dimer forms exclusively through the THJ junctions. As previously pointed out by Konarev and Popov, the energy differences between the THJ, PHHJ, and mixed THJ–PHHJ for the $[(\text{Sc}_3\text{N}@C_{80})_2]^{2-}$ dimers computed with the PBE functional and using a continuum model solvent are small.^[19] At an analogous computational level, similar results are obtained for $[(\text{Lu}_3\text{N}@C_{80})_2]^{2-}$: the three dimers are found within a range of

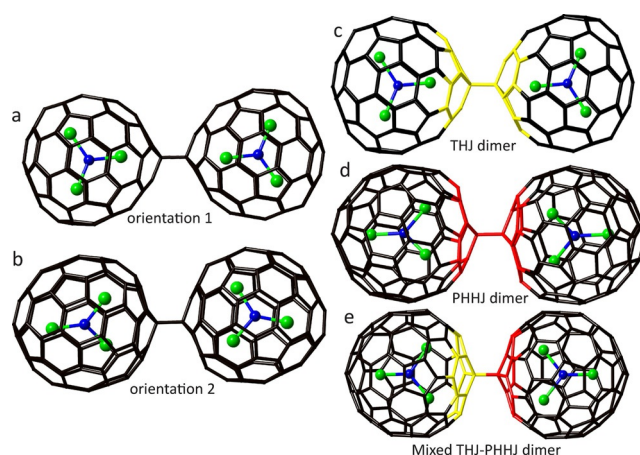


Figure 3. (a) and (b) Illustrations of the two possible M_3N cluster orientations in endohedral fullerene dimers, shown here linked through THJ–THJ bond. (c) THJ–THJ bonding with Lu_3N cluster in orientation 1 (d) PHHJ–PHHJ bonding with Lu_3N cluster in orientation 2, and (e) THJ–PHHJ mixed dimer with Lu_3N cluster in orientation 2.

only $1.1 \text{ kcal mol}^{-1}$, (see Figure 3 and Table S2). The experimentally observed THJ dimer is the lowest energy dimer for $[(\text{Lu}_3\text{N}@C_{80})_2]^{2-}$, with the mixed dimer almost at the same energy ($0.1 \text{ kcal mol}^{-1}$), followed by the symmetric PHHJ at $1.1 \text{ kcal mol}^{-1}$. These very small energy differences conflict with our experimental observation that $[\text{Ni}_{12}\text{Te}_{12}(\text{PEt}_3)_8]_2[\text{Lu}_3\text{N}@C_{80}]_2$ contains exclusively THJ–THJ dimers, hinting to additional contributions to the total free energy of the system, such as the inclusion of the zero-point energies and/or the thermal and entropic contributions.

When these contributions are considered in the calculations, the relative free energies of THJ–PHHJ and PHHJ–PHHJ dimers increase compared to the THJ–THJ which becomes somewhat more stabilized (see Table S4). We have examined the effect of reaction temperature on each type of dimer by accounting for the zero-point energies (ZPE) and the thermal and entropic contributions in the calculations within the rigid rotor and harmonic oscillator (RRHO) approximation. The general trend is that the THJ dimer is the most abundant isomer over the whole temperature range analyzed here (see Figure S1 in Supporting Information), regardless of the density functional used.

To evaluate the relevance of the stabilizing effect of the environment around the dianion we have represented the molecular electrostatic potential (MEP) distribution of the THJ dimer with and without solvent. Notice that in a continuum solvent model, both solvent and counteraction effects are included in the calculations. In both cases, the region around the inter-cage bond has the highest electron density (shown in red in Figure S2). Figure S2 shows that the electron density at the inter-cage junction increases significantly when the solvent environment is included in the calculation. These results suggest that the electrostatic environment surrounding the fullerenes can increase the stability of the dimer system by promoting the accumulation of electron density in the bonding hemispheres. A similar process could be at play in $[\text{Ni}_{12}\text{Te}_{12}(\text{PEt}_3)_8]_2[\text{Lu}_3\text{N}@C_{80}]_2$ crystal as the cluster cations are located near the nucleophilic regions.

We computed spin density distribution for the monomeric radical anion $[\text{Lu}_3\text{N@C}_{80}]^-$ (calculated for both THJ and PHHJ in orientation 1). In the THJ case, the C atom at the junction point holds the largest spin density while that in the PHHJ has a smaller spin density (Figure S3) which is more distributed on the fullerene cage. In agreement with the spin density distribution, the HOMO and the LUMO of the THJ $[\text{Lu}_3\text{N@C}_{80}]$ dimer are essentially localized on the cage, with the HOMO describing the bond formed between the two moieties (Figure 4).

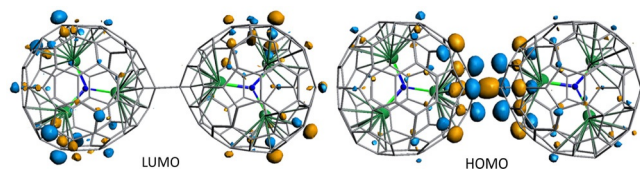


Figure 4. Representation of the LUMO and HOMO of the THJ $[\text{Lu}_3\text{N@C}_{80}]_2^{2-}$ dimer with Lu_3N in orientation 1.

Our results demonstrate that the nature of the encapsulated metal cluster controls the relative stability and orientation of the dimerization product. Konarev and Popov investigated the $(\text{Sc}_3\text{N@C}_{80})_2^{2-}$ dimer through DFT calculations and found that the energy difference between the THJ and PHHJ dimers is small (less than 2 kcal mol^{-1}), in good agreement with their experimental observation that two types of dimers are present in the crystal structure. Our calculations agree well with these results and predict that the THJ dimer becomes energetically favored for $(\text{Lu}_3\text{N@C}_{80})_2^{2-}$ and even more so for $(\text{Y}_3\text{N@C}_{80})_2^{2-}$. The crystal structure of $[\text{Ni}_{12}\text{Te}_{12}(\text{PEt}_3)_8]_2[(\text{Lu}_3\text{N@C}_{80})_2]$ reported in this communication is consistent with this first prediction.

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Conflict of interest

The authors declare no conflict of interest.

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